

# Phenomenology of Majorana zero modes in full-shell hybrid nanowires

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May 30, 2024



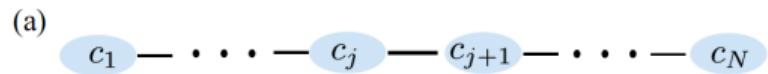
# Outline

- ① Engineering topologically protected edge states
- ② Signals in the LDOS: CdGM analogs
- ③ Opening the topological minigap
- ④ Conclusions

# The Kitaev chain

- ▶ Chain of  $N$  spin-less fermions ( $p$ -wave superconductivity):

$$H = -\mu \sum_{j=1}^N \left( c_j^\dagger c_j - \frac{1}{2} \right) + \sum_{j=1}^{N-1} \left[ -t (c_j^\dagger c_{j+1} + c_{j+1}^\dagger c_j) + \Delta (c_j c_{j+1} + c_{j+1}^\dagger c_j^\dagger) \right]$$



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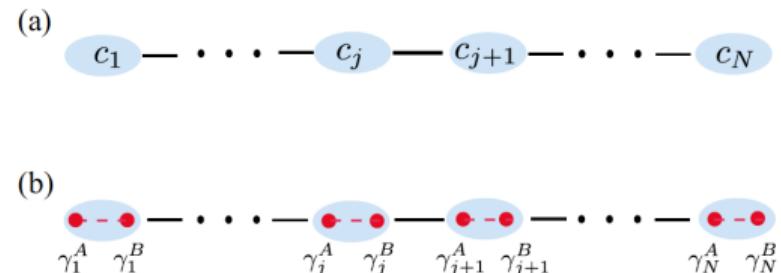
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- Majorana representation:

$$c_j = \frac{1}{2} (\gamma_j^A + i\gamma_j^B), \quad c_j^\dagger = \frac{1}{2} (\gamma_j^A - i\gamma_j^B)$$



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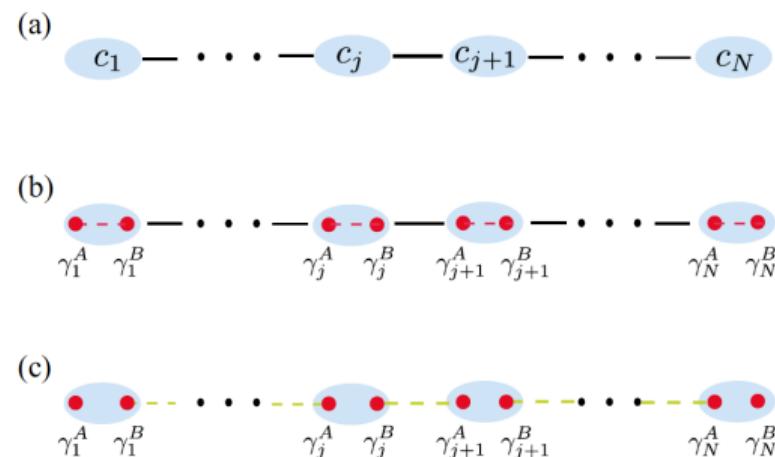
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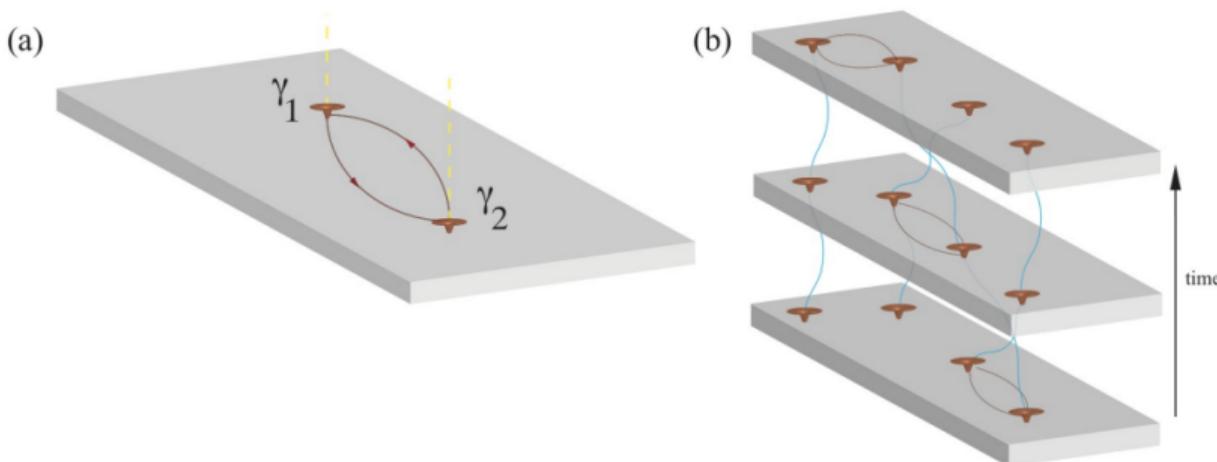
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- Hamiltonian in terms of Majorana operators:

$$H = -\frac{i\mu}{2} \sum_{j=1}^N \gamma_j^A \gamma_j^B + \frac{i}{2} \sum_{j=1}^{N-1} [(\Delta + t) \gamma_j^B \gamma_{j+1}^A + (\Delta - t) \gamma_j^A \gamma_{j+1}^B]$$



# Majoranas for qubits



- ▶ MZM are non-Abelian anyons.
- ▶ Gap closing/reopening  $\Rightarrow$  topological protection.

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# We need a $p$ -wave superconductor!

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$$\Delta \left( c_j c_{j+1} + c_{j+1}^\dagger c_j^\dagger \right).$$

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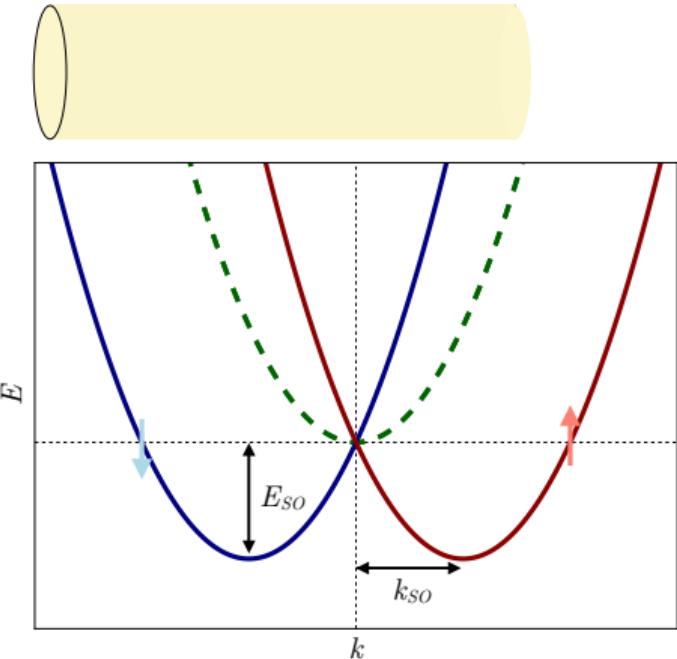
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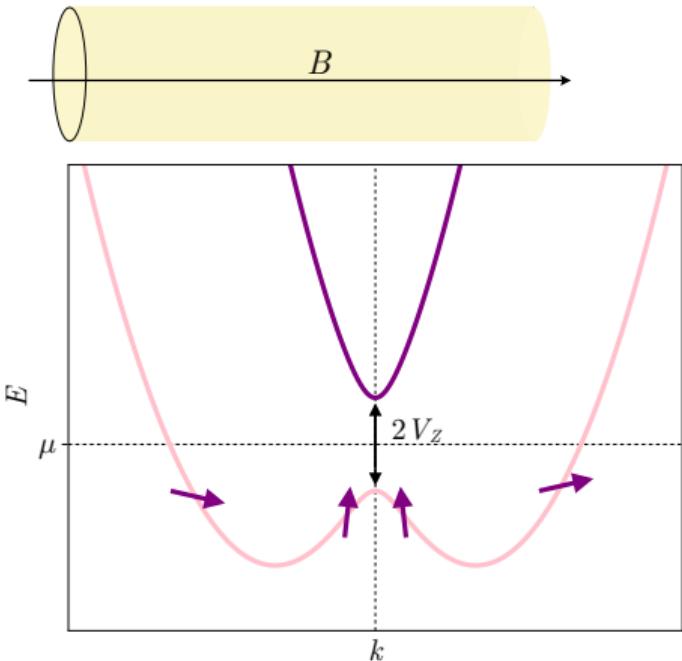
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- ▶ Lutchyn and Oreg: proximitize semiconductors with strong spin-orbit coupling.

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# Rashba, Zeeman and helical bands

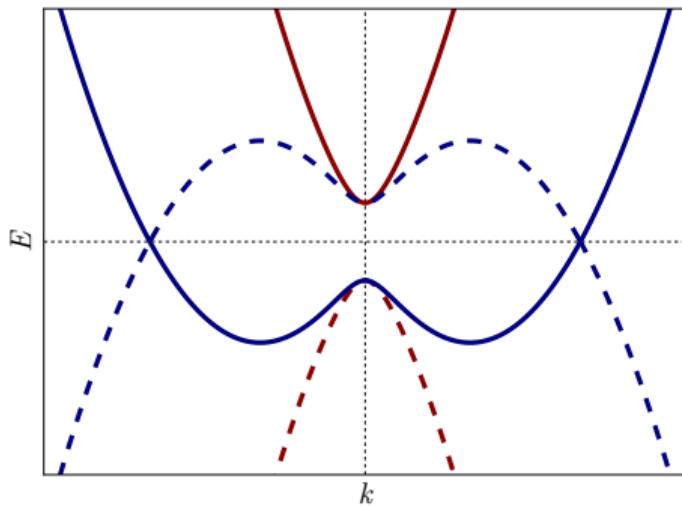
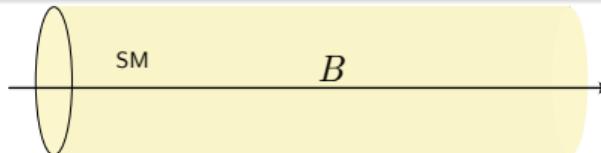
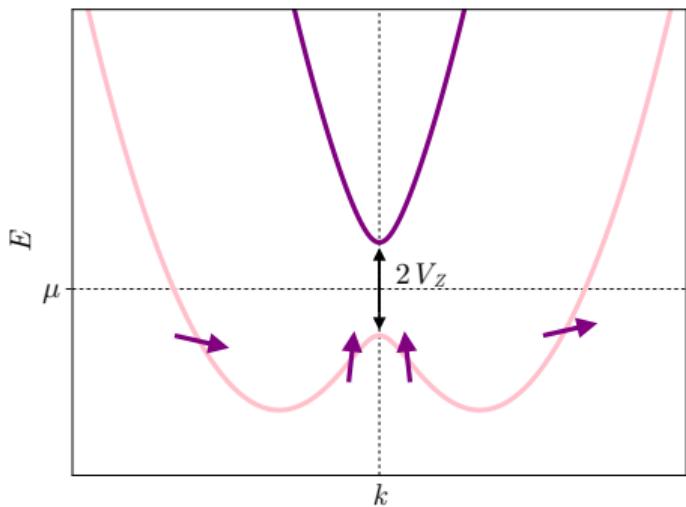
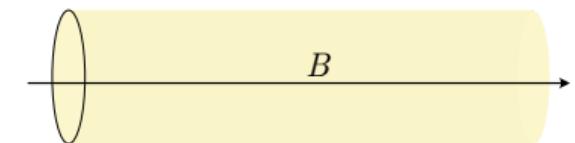


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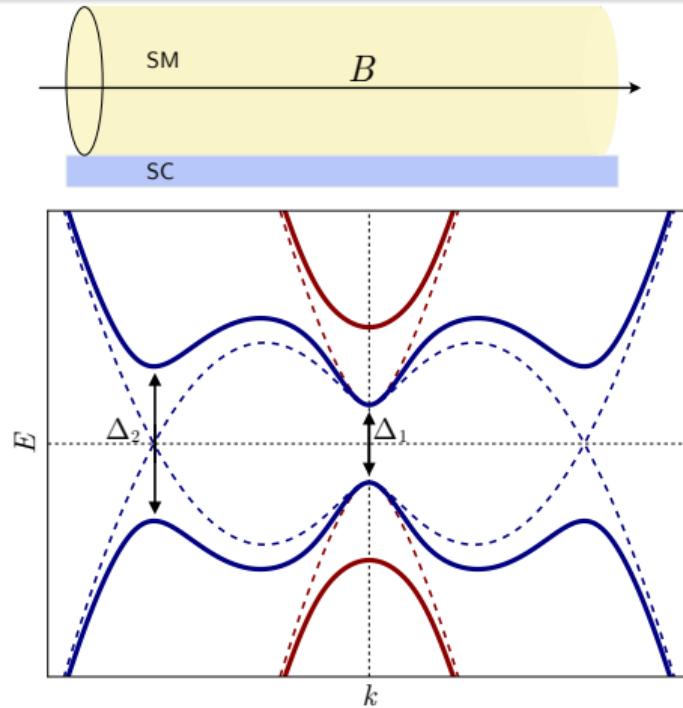
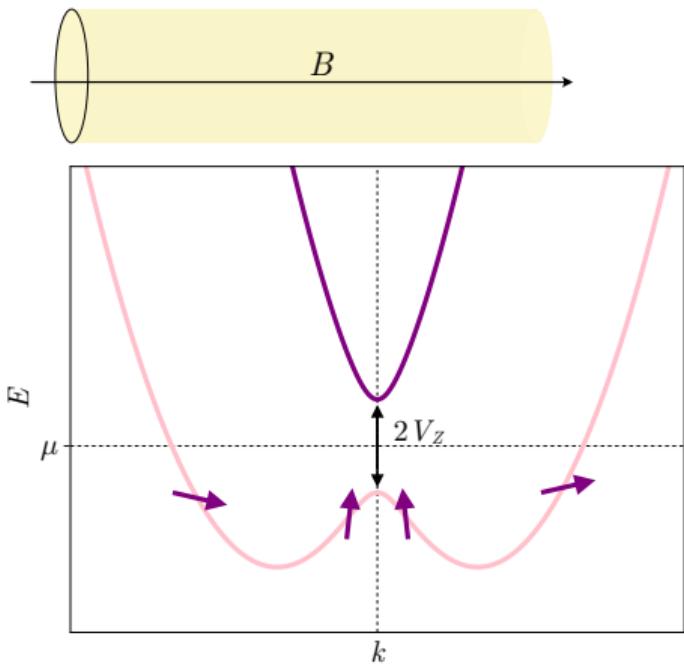
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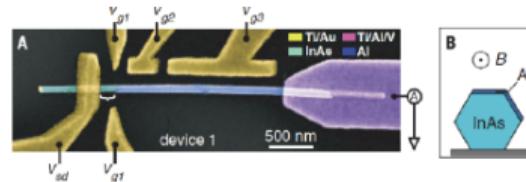
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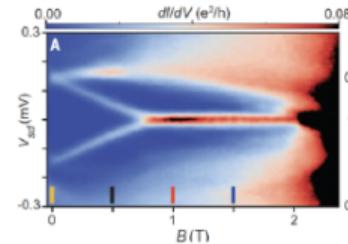
# Searching for Majoranas

- ▶ Strong experimental interest.



Zero-bias anomalies in tunneling spectroscopy experiments

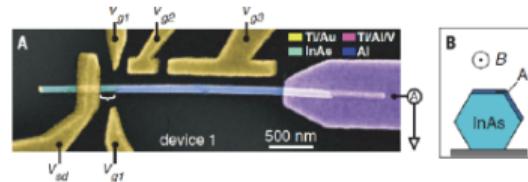
Mourik 2012, Albrecht 2016, Deng 2016



Claims: V. Mourik *et al.* 2012, *Science*. S. M. Albrecht *et al.* 2016, *Nature*. M. T. Deng *et al.* 2016, *Science*.  
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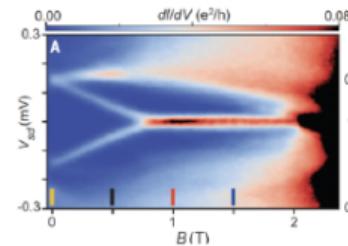
# Searching for Majoranas

- ▶ Strong experimental interest.
- ▶ Zero-bias anomalies detected with non-topological explanations.



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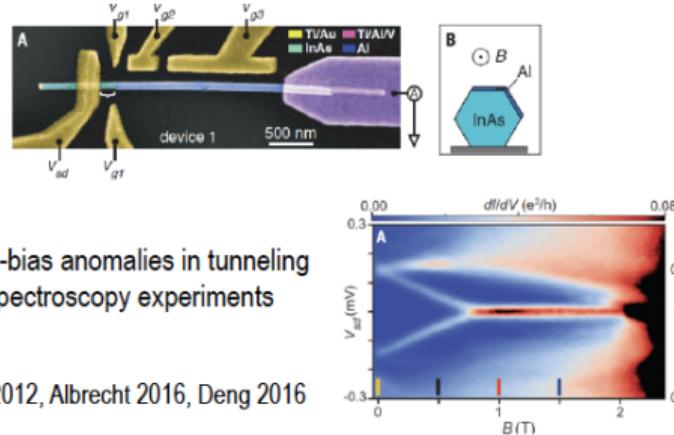
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- ▶ Drawbacks:

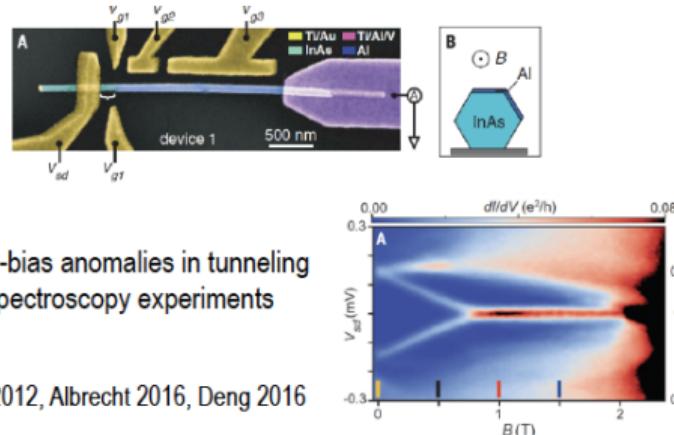


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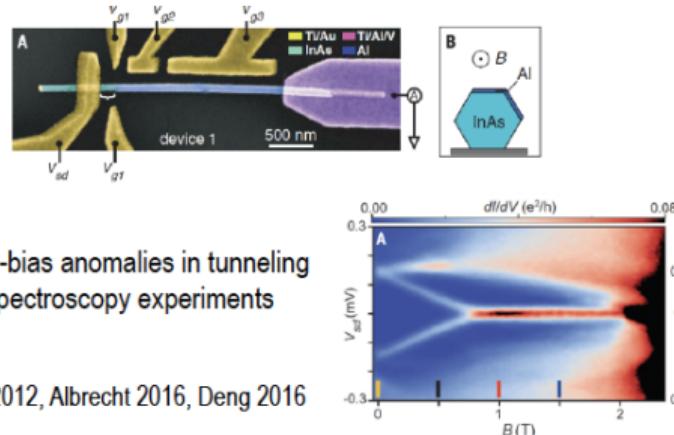
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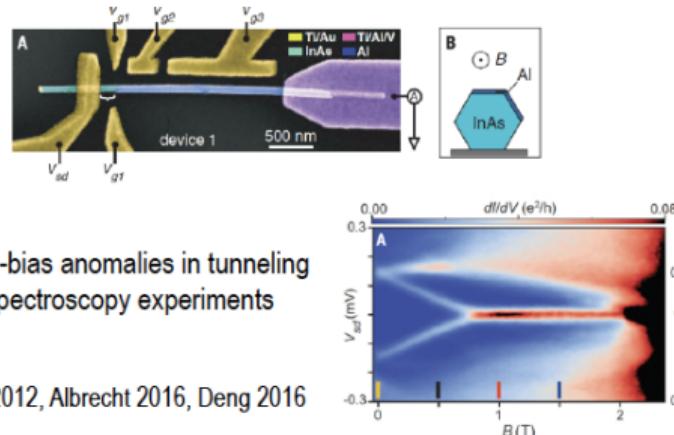
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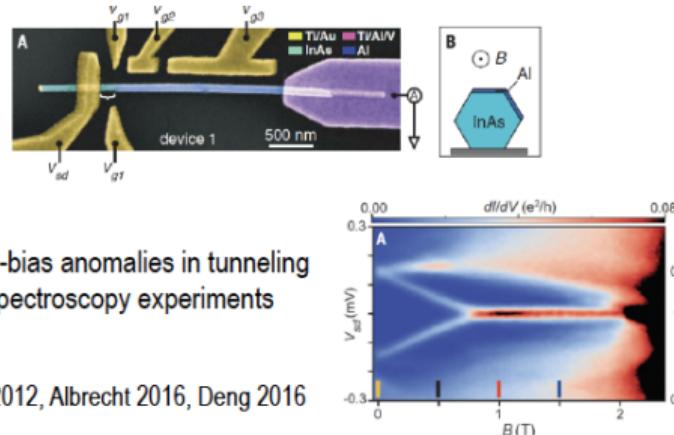
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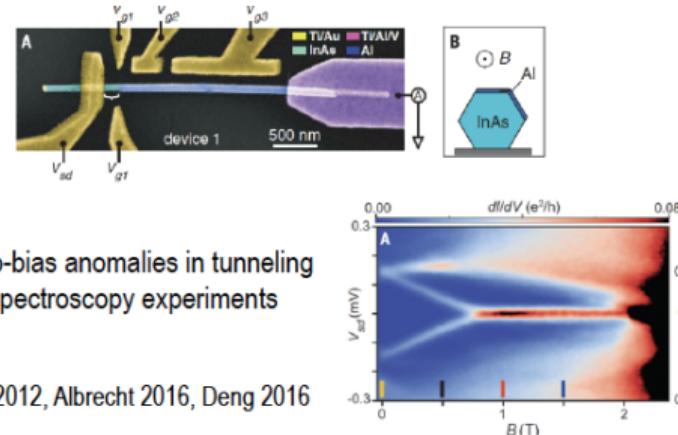


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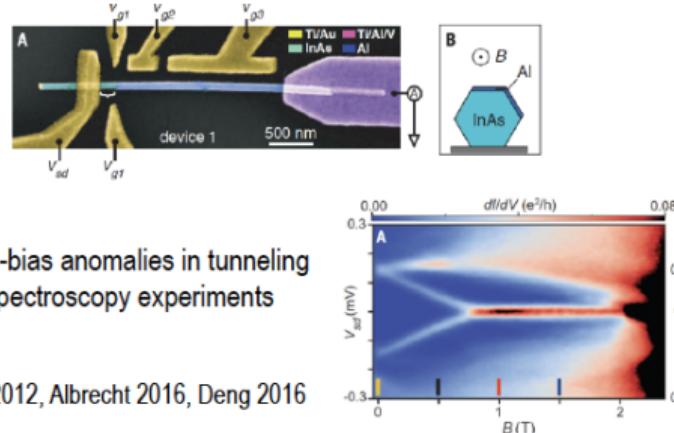


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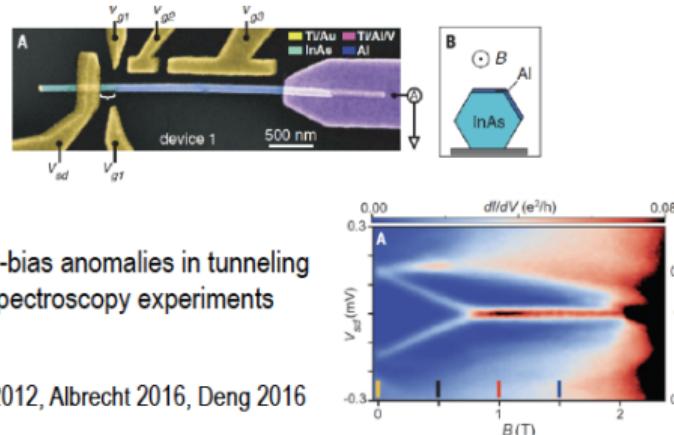
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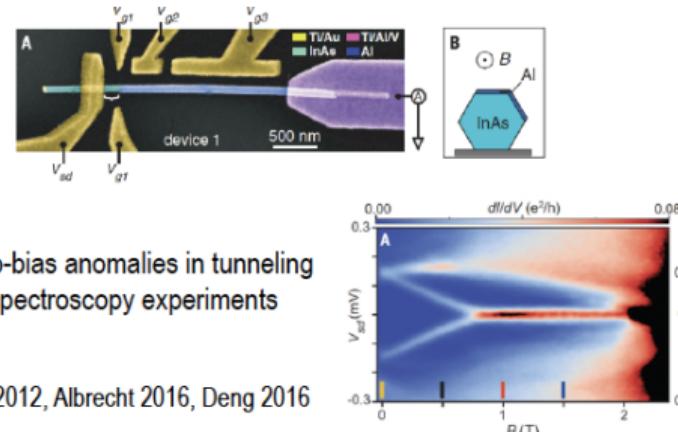
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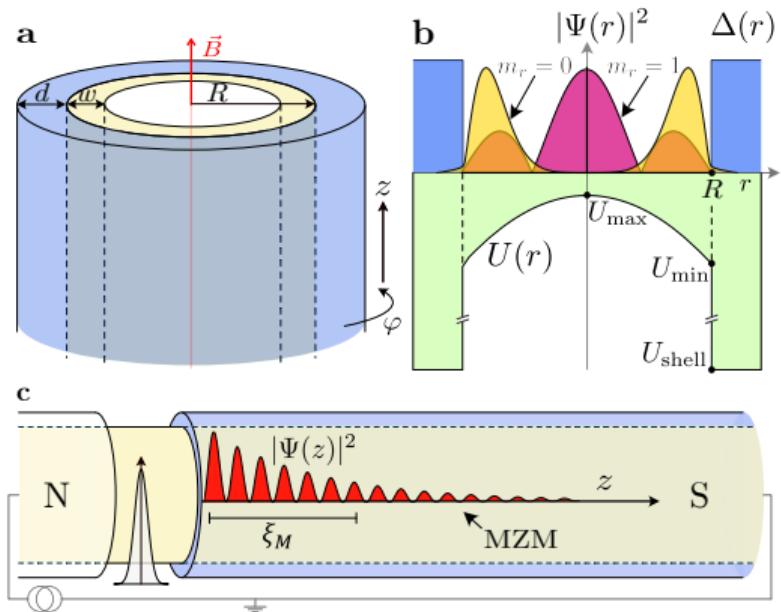
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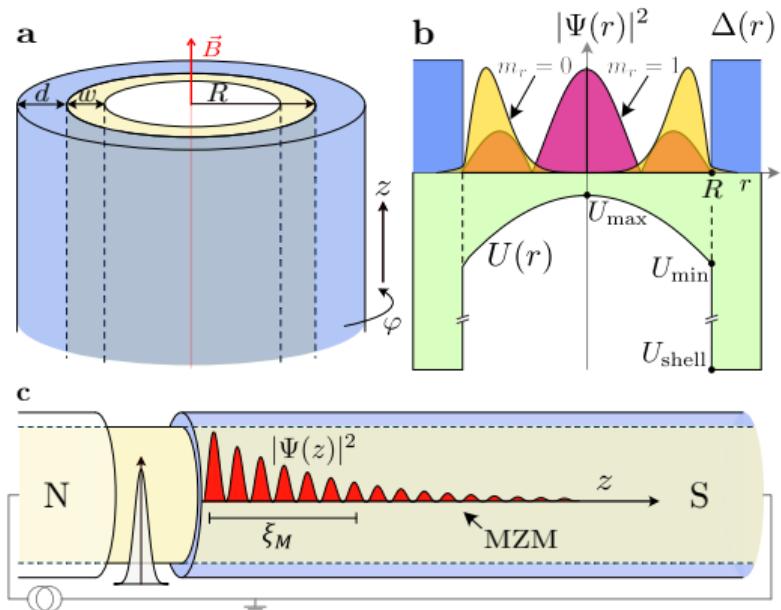
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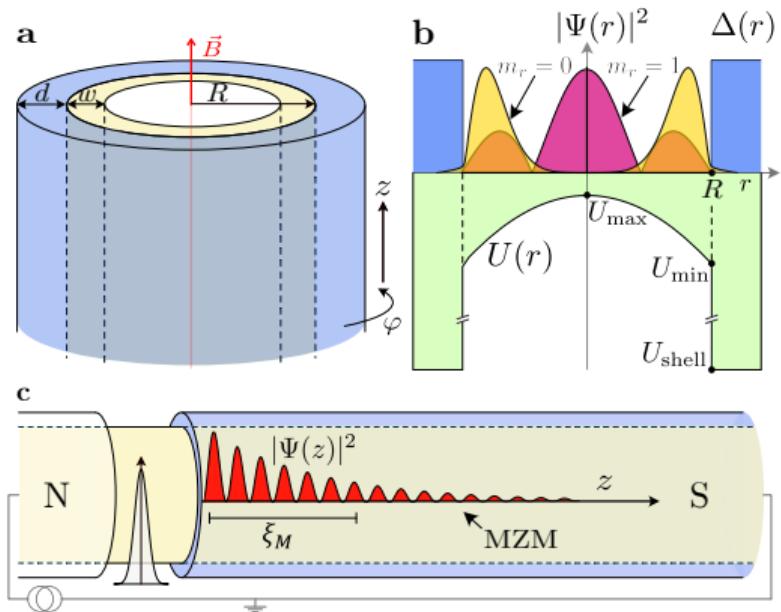
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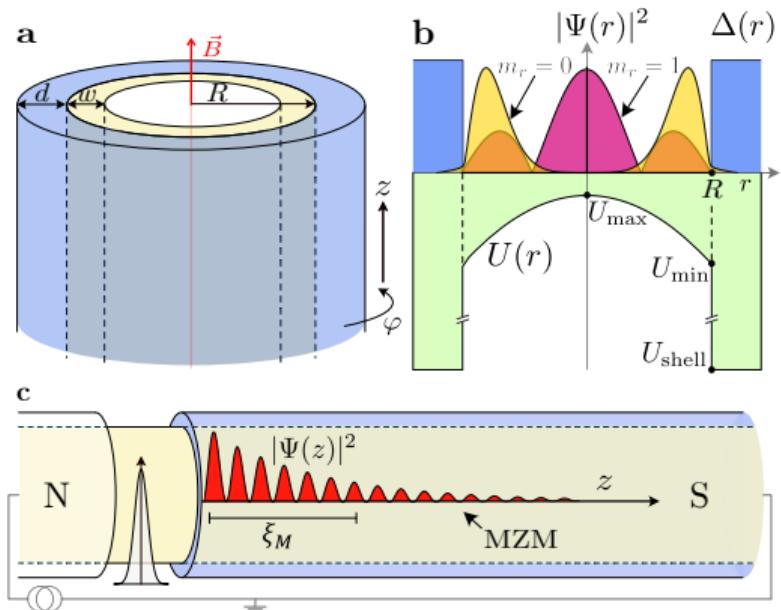
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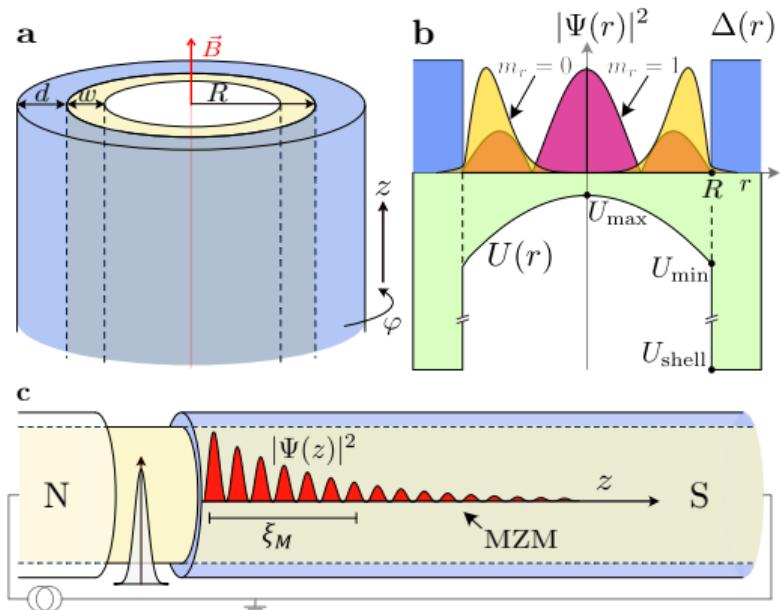
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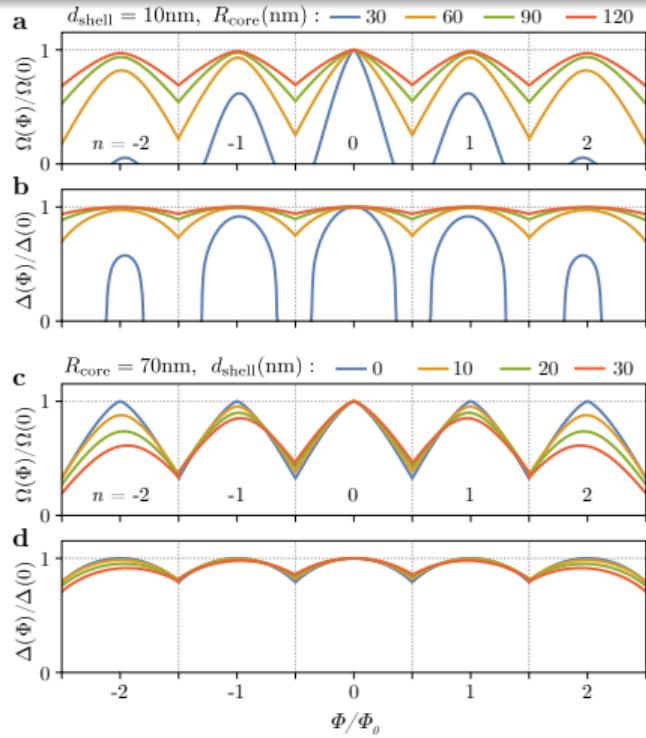


## ► Key points:

- Cylindrical symmetry
- Topological transition driven by orbital effect **not Zeeman**
- Needs lower magnetic fields.
- Only one angular mode can be topological.

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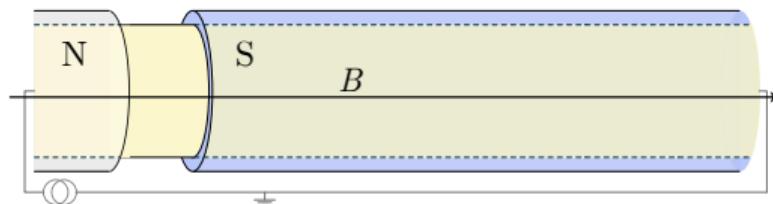
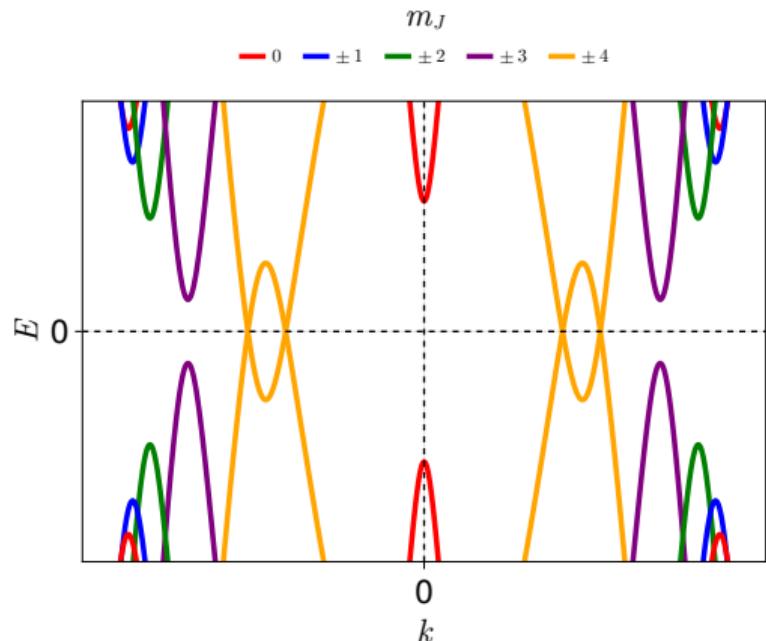
# The Little-Parks effect



- ▶ Cylinder  $\Leftrightarrow$  vortex.
- ▶ Too thin for full Meissner.
- ▶ Quantized winding of the order parameter:  $\Delta = |\Delta| e^{in\varphi}$ .
- ▶  $n \in \mathbb{Z}$  and jumps every flux quantum  $\Phi_0$ .
- ▶ Quasi-quantization of flux  $\Rightarrow$  pairing presents LP lobes.
- ▶ Depends on  $R$ , SC thickness  $d$  and  $\xi_d$ , the SC coherence length.

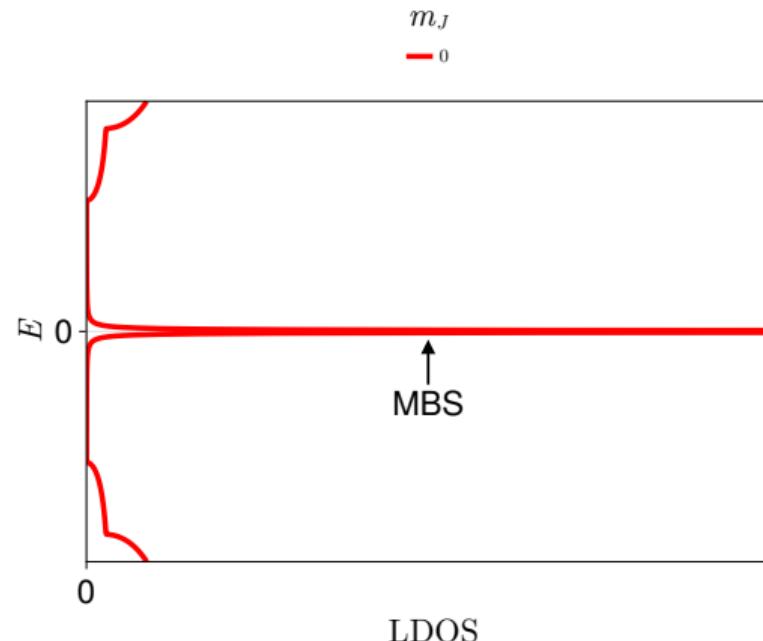
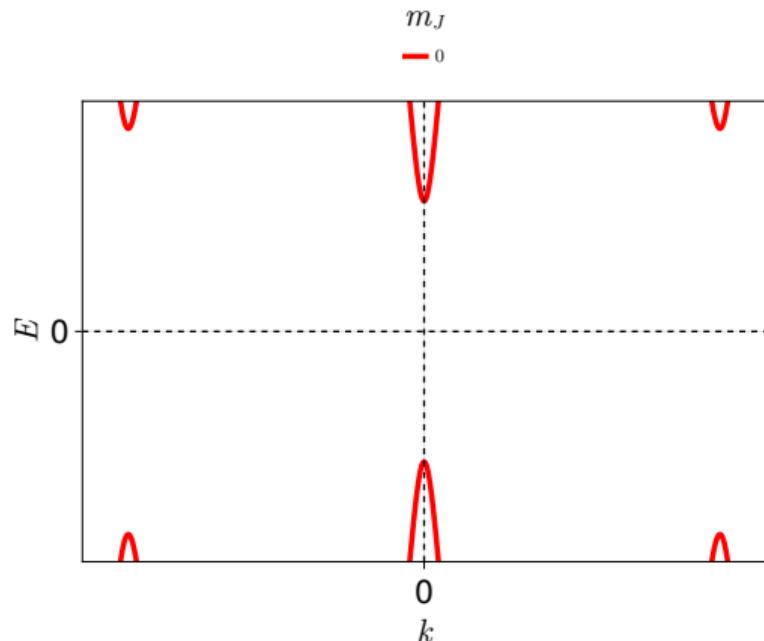
W. A. Little and R. D. Parks 1962, *Phys. Rev. Lett.*  
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# The CdGM analog states



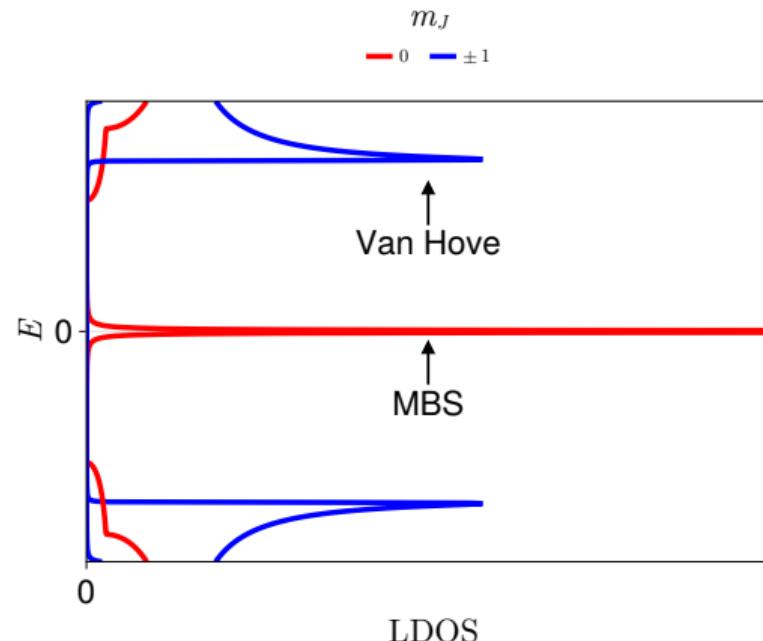
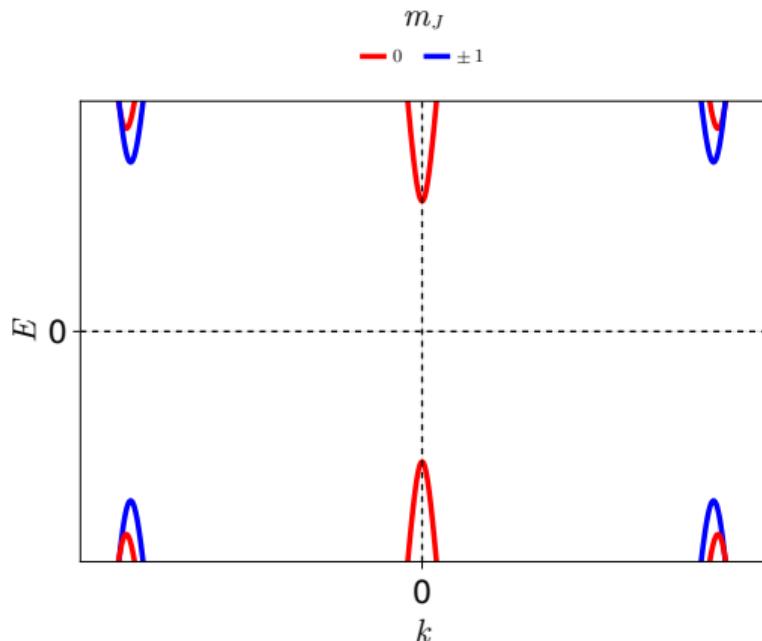
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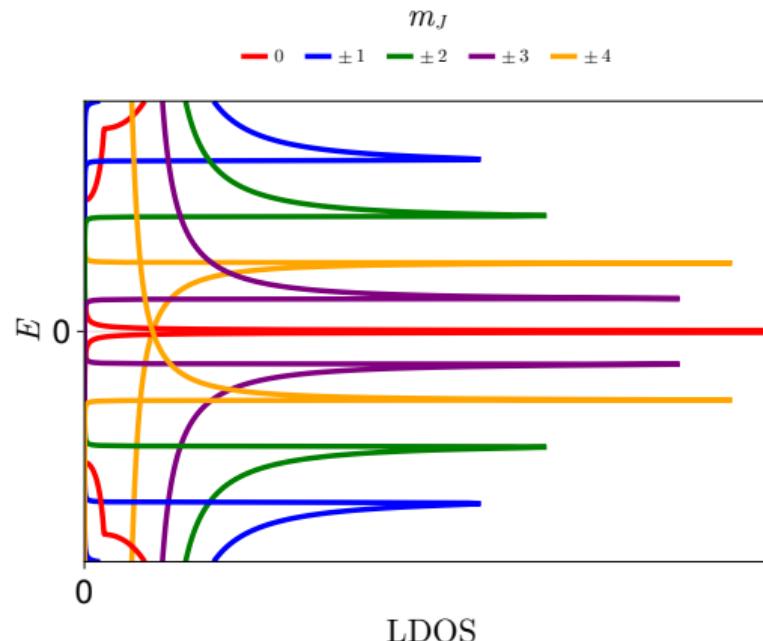
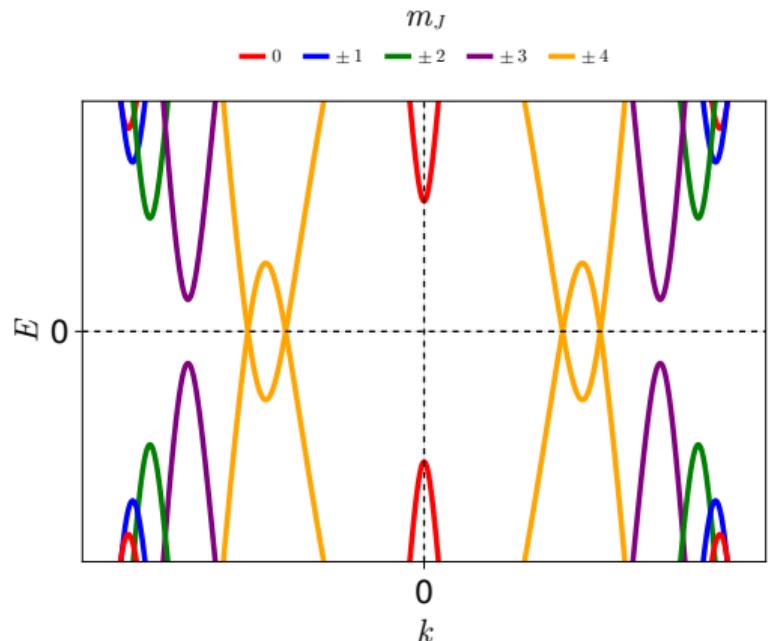
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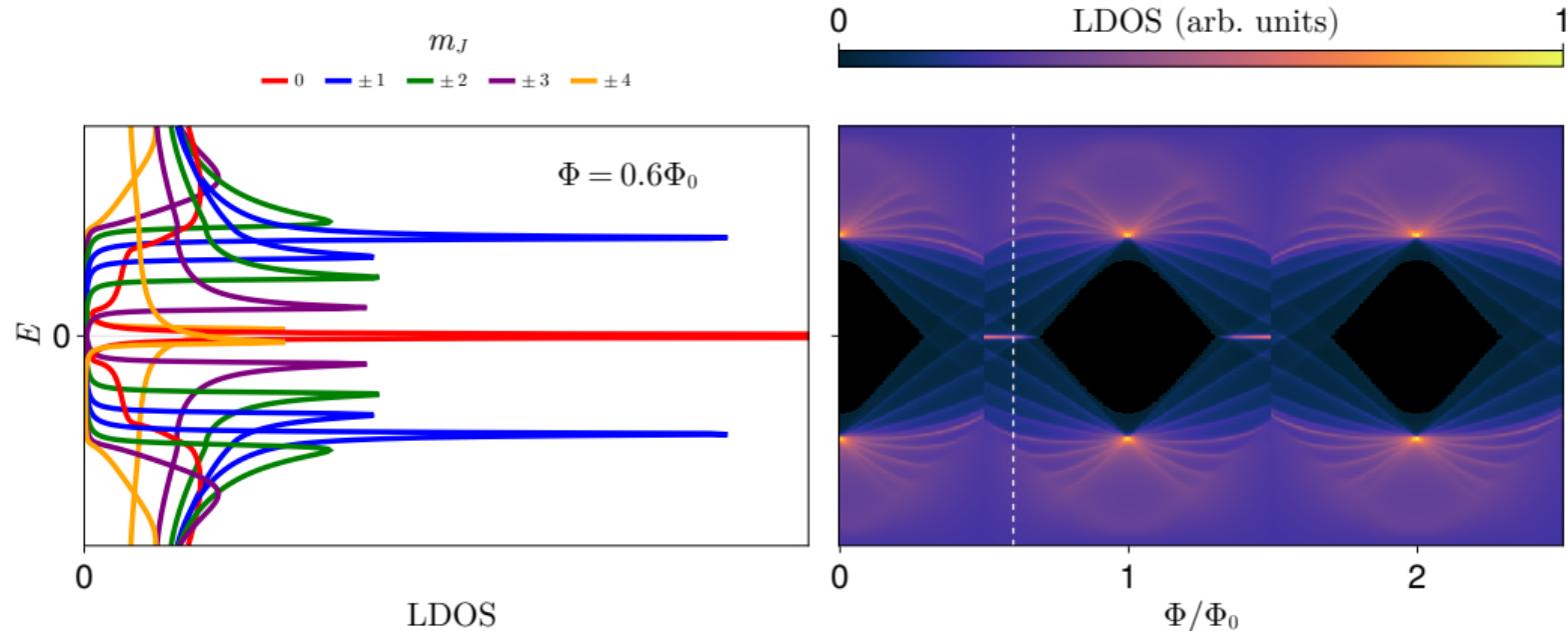
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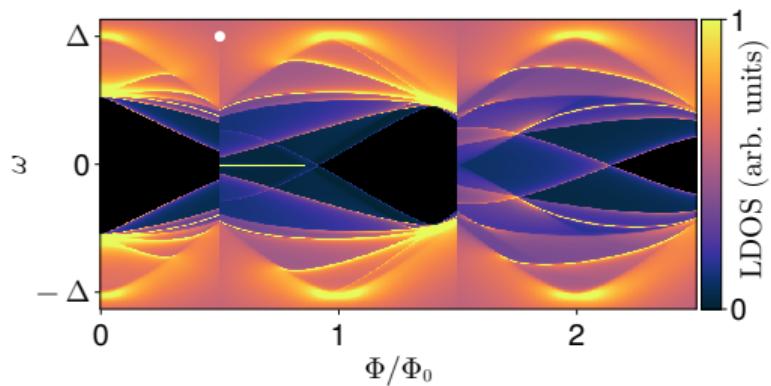
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# LDOS vs. flux



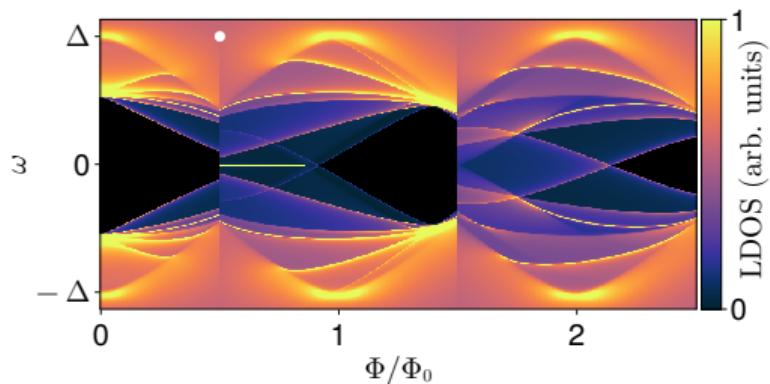
C. Payá et al. 2024, Phys. Rev. B.  
P. San-Jose et al. 2023, Phys. Rev. B.

# The tubular-core model



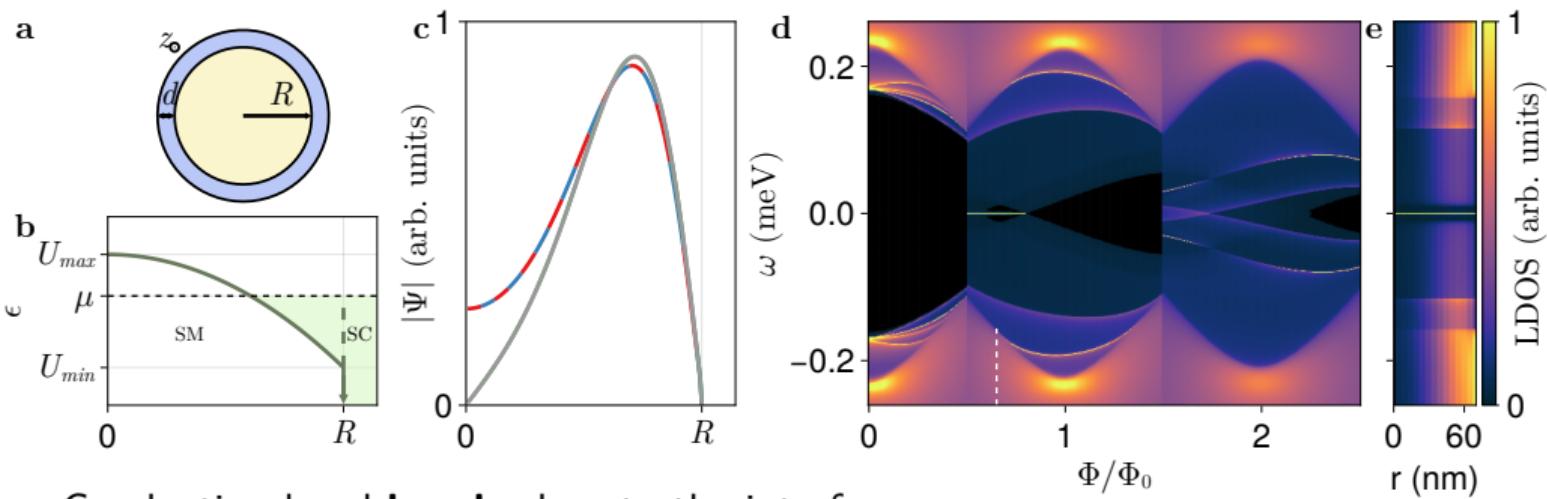
- ▶ Adding a width to the semiconductor.

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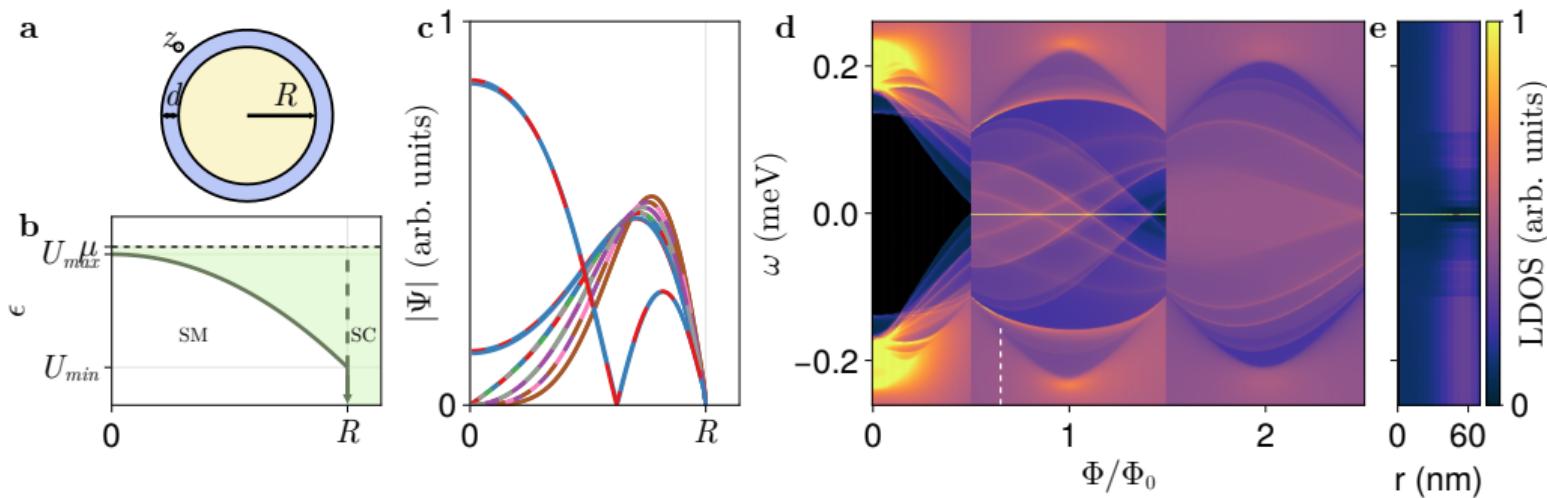
- ▶ Adding a width to the semiconductor.
- ▶ Most common scenario: CdGMs fill the MZM minigap.
- ▶ No topological protection

# Pushing the WF to the interface

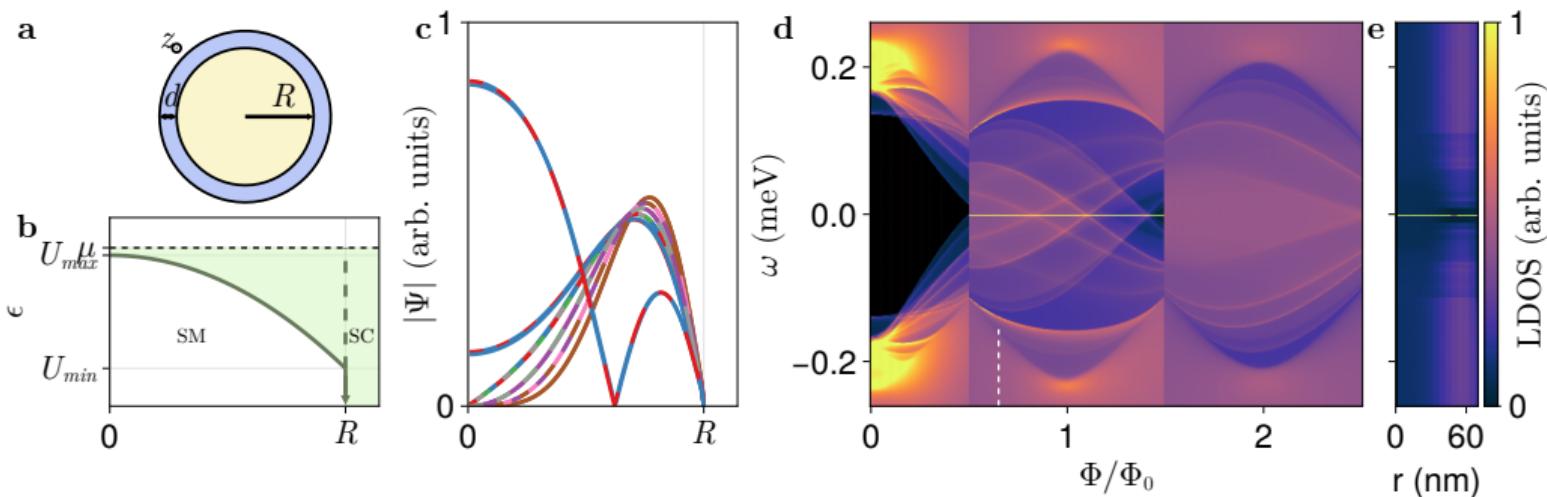


- ▶ Conduction band **bends** close to the interface.

## Second radial mode: protection lost

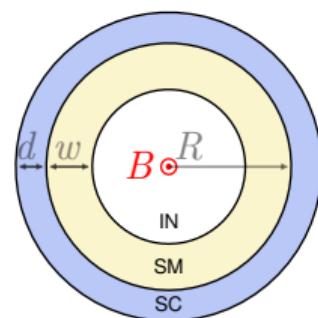
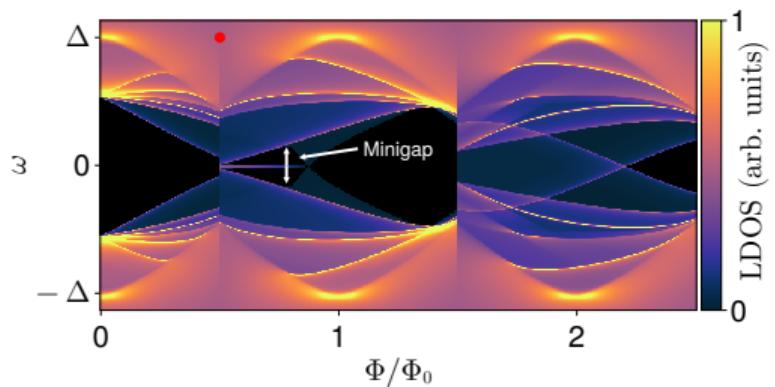


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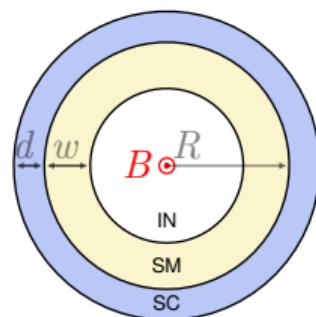
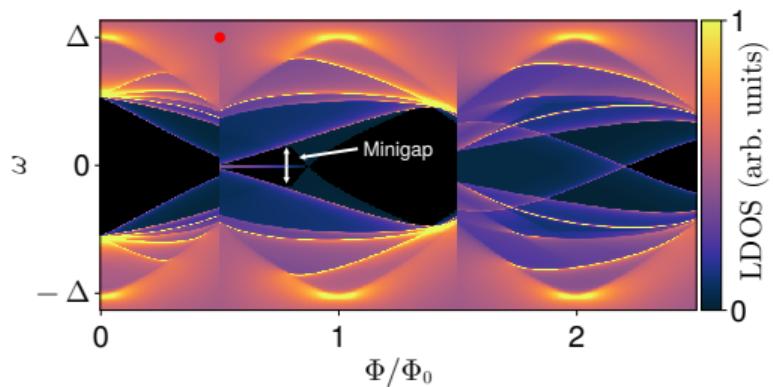
- When the second radial mode is occupied, the ZEP expands over the full lobe, but CdGMs cover it.

# Protected islands in the tubular-core



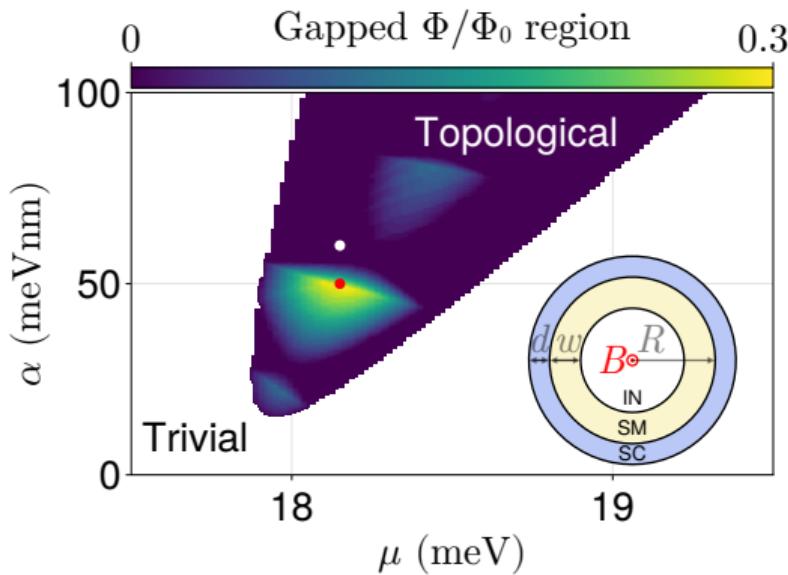
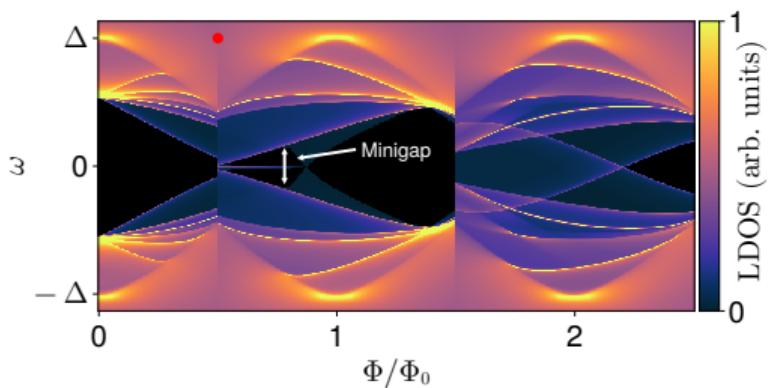
C. Payá *et al.* 2024, *Phys. Rev. B*.

## Protected islands in the tubular-core



- We need to push the charge to the interface.

## Protected islands in the tubular-core



- ▶ We need to push the charge to the interface.
- ▶ Topologically protected islands appear-

## Summary

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## Summary

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### Take home message

*Majorana physics of full-shell nanowires is very rich. For pristine configurations, the tubular-core model is the optimal candidate in comparison to the solid-core geometry.*

# People involved

Project Leader



Elsa Prada



Pablo San José



Ramón Aguado



## Theory:

Samuel D. Escribano (Weizmann Institute)

Andrea Vezzosi (U. of Modena, now U. of Lausanne)

Fernando Peñaranda (DIPC)

## Experimentalists:

Saulius Vaitiekėnas (Niels Bohr Institute)

Charles M. Marcus (NBI, now U. of Washington)

## Ongoing experiments:

Jesper Nygård (Niels Bohr Institute)



# Phenomenology of Majorana zero modes in full-shell hybrid nanowires

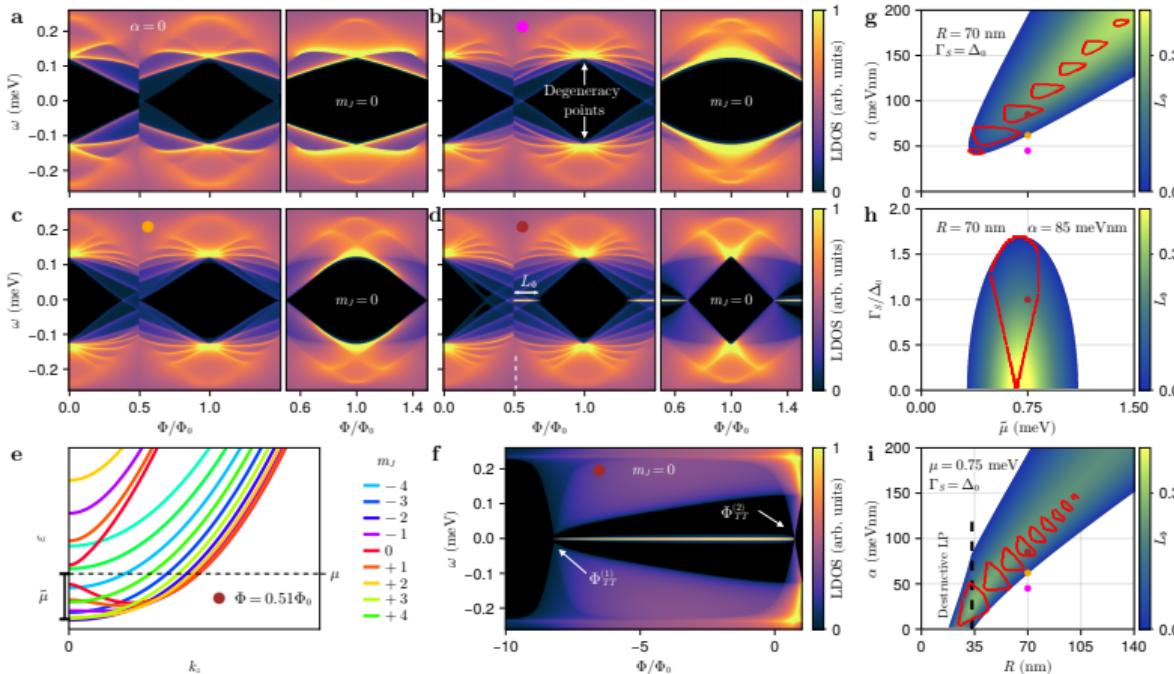
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May 30, 2024

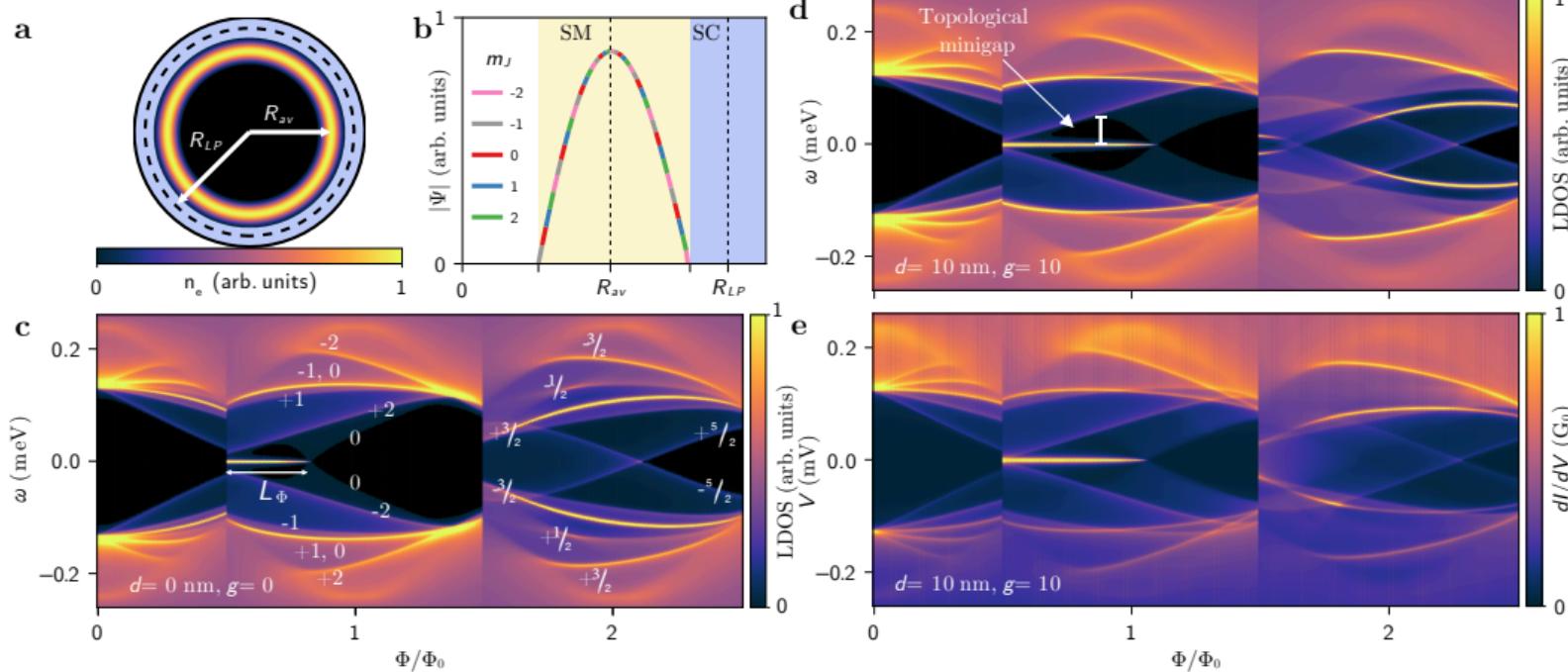


# Hollow-core results



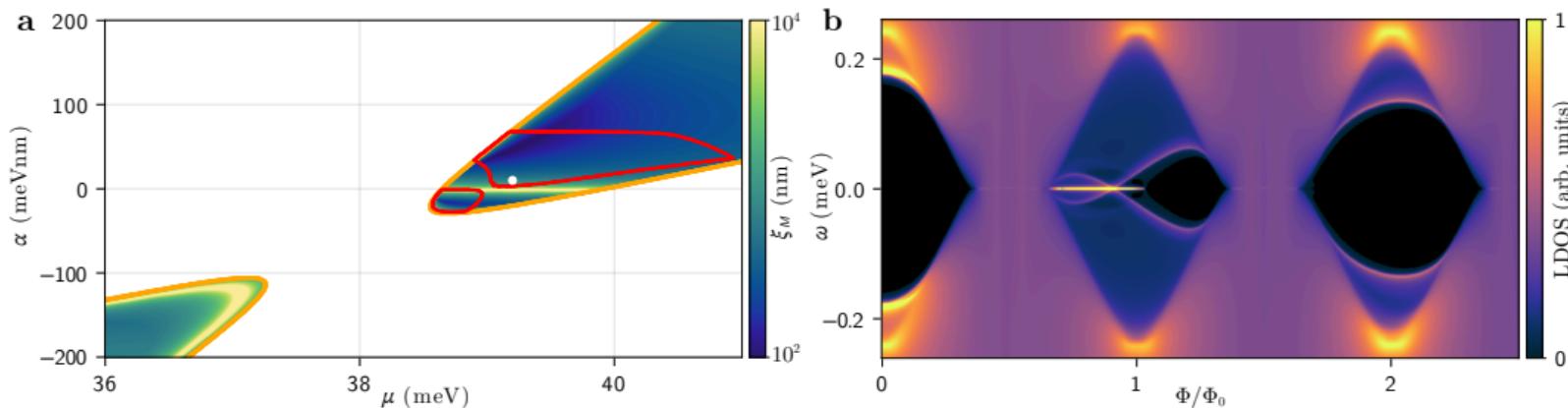
C. Payá et al. 2024, Phys. Rev. B.

# Modified hollow-core results

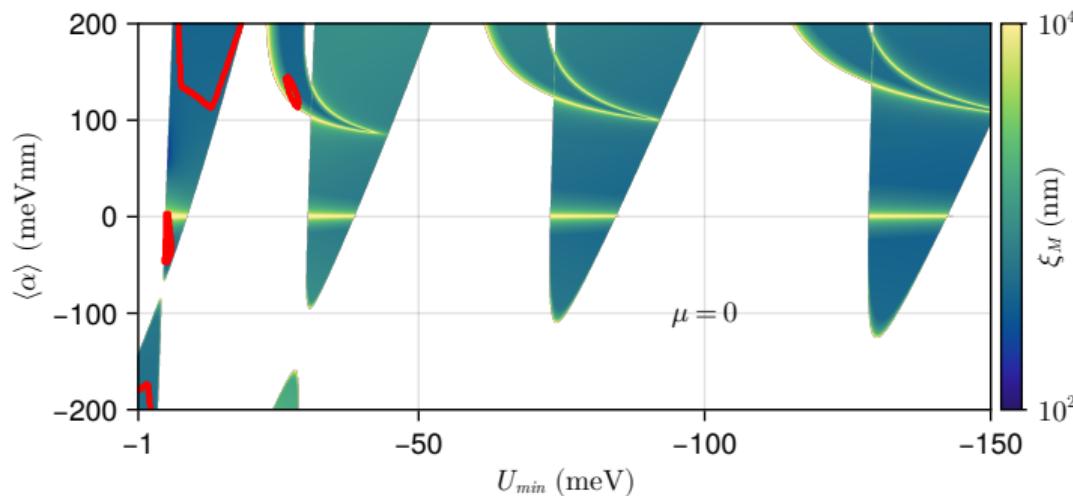


C. Payá *et al.* 2024, Phys. Rev. B.

# Destructive Little-Parks

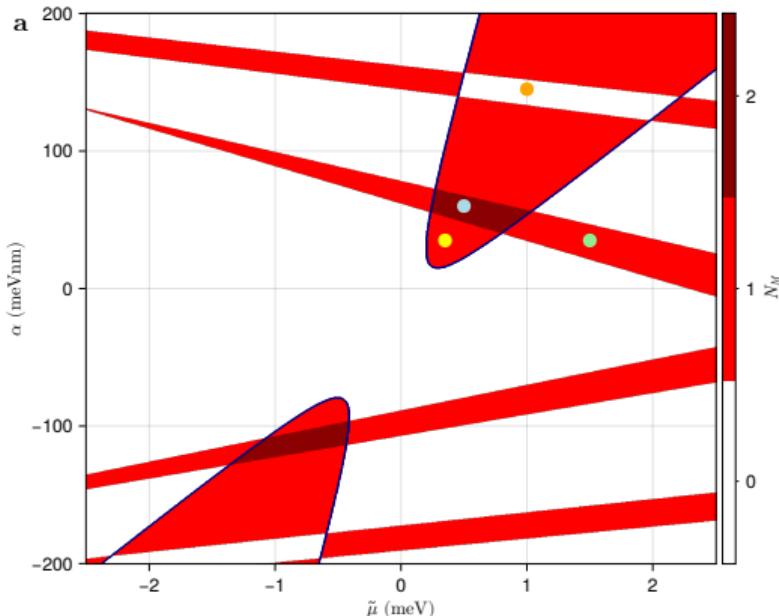


## Band-bending: not enough islands



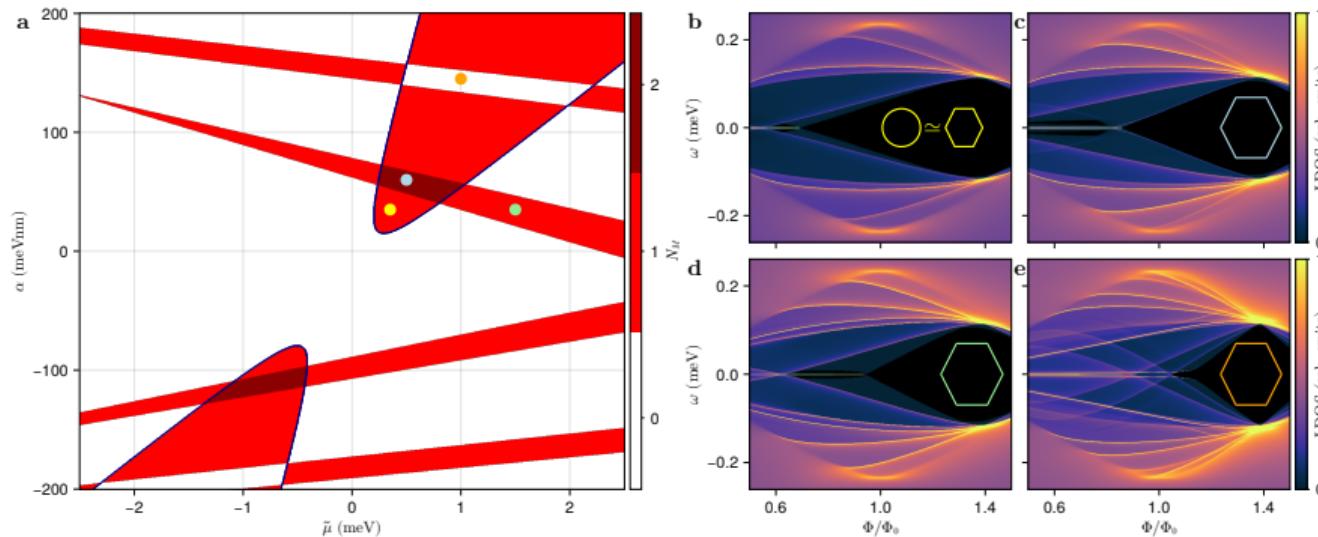
- ▶ Notice axes are mean  $\alpha$  and  $U_{min}$ , the minimum of the dome-profile.
- ▶ One wedge per radial mode. No islands outside the first radial mode.

# Hexagonal wave-function



- ▶ New red stripes. Hexagon has  $\ell = 6$ .
- ▶ Upper stripe:  $m_J = 0$  mixes with  $m_J = \pm 6$ .
- ▶ Lower stripe:  $m_J = 3$  mixes with  $m_J = -3$ .
- ▶ The MZM coming from  $m_J = \pm 3$  **cannot** interact with  $m_J = 0 \Rightarrow$  they overlap.
- ▶ The  $m_J = \pm 6$  MZM annihilates the  $m_J = 0$  MZM.

# Hexagonal wave-function



- Except for the new topological stripes and a region where the MZM splits, the system is equivalent to the cylinder.

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# Phenomenology of Majorana zero modes in full-shell hybrid nanowires

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